

## Large deflection analysis of point supported super-elliptical plates

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*(Received August 23, 2013, Revised April 3, 2014, Accepted May 18, 2014)*

**Abstract.** Nonlinear bending of super-elliptical plates of uniform thickness under uniform transverse pressure was investigated by the Ritz method. The material was assumed to be homogeneous and isotropic. The contribution of the boundary conditions at the point supports was introduced by the Lagrange multipliers. The solution was obtained by the Newton-Raphson method. The influence of the location of the point supports on the central deflection was highlighted by sensitivity analysis. An approximate relationship between the central deflection and the super-elliptical power was obtained using the method of least squares. The critical points where the maximum deflection may develop, and the influence of nonlinearity were highlighted. The nonlinearity was found to be sensitive to the aspect ratio. The accuracy of the algorithm was validated by comparing the central deflection with the solutions of elliptical and rectangular plates.

**Keywords:** super-elliptical; plate; nonlinear; bending; large deflection

### 1. Introduction

If the deflection of a plate is beyond a certain level (i.e.,  $w \geq 0.3t$ ) the relation between external load and deflection is no longer linear (Szilard 2004). It is well known that when the deflections are of the same order as the thickness, the membrane forces play a more significant role in carrying the load (Alwar and Nath 1976, Gorji and Akileh 1990), and thus, the analysis must be extended to include the additional effects produced by large deflections (Szilard 2004).

The nonlinear analysis of plates has been a challenging topic in applied mechanics due to the coupled governing differential equations. Since closed form solutions are only available for a limited number of cases depending on the geometry, boundary conditions, material properties, and loading, the problem involving nonlinearity has been attacked by means of various numerical methods (Silverman and Mays 1972, Civalet 2005, Malekzadeh 2007, Dai *et al.* 2013).

Plates of various shapes have been studied by different methods using several plate theories (Mukhopadhyay and Bera 1994, Bayer *et al.* 2002, Özkul and Türe 2004, Artan and Lehmann 2009, Orakdöğen *et al.* 2010, Kutlu and Omurtag 2012, Chen 2013). As far as the author knows, in the literature devoted to the nonlinear analysis of plates, except the only paper in which clamped and simply supported boundaries were examined (Zhang 2013), no study has been published on the large deflection of super-elliptical plates. Super-elliptical plates include a large variety of plate

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shapes ranging from an ellipse to a rectangle with rounded corners. Rectangular plates with rounded corners enable to diffuse and dilute stress concentrations unlike plates with sharp corners (Liew *et al.* 1998). However, although they are extensively used as structural and machine elements in aviation, shipbuilding, bridge construction and in applications of aerospace and naval engineering where the deflections are no longer small (Zhang and Kim 2006), the number of studies dealing with them has been rather limited (Asemi *et al.* 2013, Ceribasi 2013, Hasheminehad *et al.* 2013, Jazi and Farhatnia 2012, Tang *et al.* 2012, Algazin 2011, Wu and Liu 2005, Zhou *et al.* 2004, Pedersen 2004, Liew and Feng 2001, Chen *et al.* 1999, Lim *et al.* 1998, Wang *et al.* 1994, and the references cited therein). Besides, plates resting on isolated points such as telescope mirrors, solar panels, printed circuit boards, or slabs supported by columns are frequently encountered in structural design. In the current study the nonlinear bending behavior of a super-elliptical plate undergoing large deflection, and supported symmetrically by four intermediate point supports on the diagonals was examined. Sensitivity analysis was made to determine the influence of the position of the point supports on the maximum deflection. Since the bending moments are crucial in design, the support location minimizing the bending moments at the supports and at the center of the plate was investigated. The large deflection bending response of the plate was compared with the linear solution, and the influence of the thickness on the central deflection was observed. Some of the numerical results were presented for future reference. The accuracy of the algorithm was checked with the solutions of elliptical and rectangular plates, and good agreement was obtained.

## 2. Formulation

In Cartesian coordinates the boundary contour of a super-elliptical plate can be defined by (Liew *et al.* 1998)

$$\left(\frac{x}{a}\right)^{2k} + \left(\frac{y}{b}\right)^{2k} = 1, \quad k = 1, 2, \dots, \infty \quad (1)$$

The strain energy of the plate due to bending, the strain energy of the plate due to strain of the middle surface, the potential energy of the lateral load, and the bending moments can be written as (Timoshenko and Woinowsky-Krieger 1959)

$$V_1 = \frac{D}{2} \int_{x_1}^{x_2} \int_{y_1}^{y_2} \left\{ \left[ \left( \frac{\partial^2 w}{\partial x^2} \right) + \left( \frac{\partial^2 w}{\partial y^2} \right) \right]^2 - 2(1-\mu) \left[ \left( \frac{\partial^2 w}{\partial x^2} \right) \left( \frac{\partial^2 w}{\partial y^2} \right) - \left( \frac{\partial^2 w}{\partial x \partial y} \right)^2 \right] \right\} dx dy \quad (2)$$

$$V_2 = \frac{Et}{2(1-\mu^2)} \int_{x_1}^{x_2} \int_{y_1}^{y_2} \left\{ \left[ \left( \frac{\partial u}{\partial x} \right)^2 + \left( \frac{\partial u}{\partial x} \right) \left( \frac{\partial w}{\partial x} \right)^2 + \left( \frac{\partial v}{\partial y} \right)^2 + \left( \frac{\partial v}{\partial y} \right) \left( \frac{\partial w}{\partial y} \right)^2 \right] + \frac{1}{4} \left[ \left( \frac{\partial w}{\partial x} \right)^2 + \left( \frac{\partial w}{\partial y} \right)^2 \right]^2 + 2\mu \left[ \left( \frac{\partial u}{\partial x} \right) \left( \frac{\partial v}{\partial y} \right) + \frac{1}{2} \left( \frac{\partial v}{\partial y} \right) \left( \frac{\partial w}{\partial x} \right)^2 + \frac{1}{2} \left( \frac{\partial u}{\partial x} \right) \left( \frac{\partial w}{\partial y} \right)^2 \right] \right\} dx dy$$

$$\left[ + \frac{(1-\mu)}{2} \left[ \left( \frac{\partial u}{\partial y} \right)^2 + 2 \left( \frac{\partial u}{\partial y} \right) \left( \frac{\partial v}{\partial x} \right) + \left( \frac{\partial v}{\partial x} \right)^2 \right. \right. \\ \left. \left. + 2 \left( \frac{\partial u}{\partial y} \right) \left( \frac{\partial w}{\partial x} \right) \left( \frac{\partial w}{\partial y} \right) + 2 \left( \frac{\partial v}{\partial x} \right) \left( \frac{\partial w}{\partial x} \right) \left( \frac{\partial w}{\partial y} \right) \right] \right] \quad (3)$$

$$V_3 = - \int_{x_1}^{x_2} \int_{y_1}^{y_2} q w dx dy, \quad m_x = -D \left( \frac{\partial^2 w}{\partial x^2} + \mu \frac{\partial^2 w}{\partial y^2} \right), \quad m_y = -D \left( \frac{\partial^2 w}{\partial y^2} + \mu \frac{\partial^2 w}{\partial x^2} \right) \quad (4)$$

$$x_1 = -x_2 = -a, \quad y_1 = -y_2 = -\frac{b}{a} \sqrt[2k]{a^{2k} - x^{2k}}, \quad D = \frac{Et^3}{12(1-\mu^2)} \quad (5)$$

$$W = \frac{w}{t} = \frac{qa^4 \alpha}{Dt}, \quad U = \frac{u}{t}, \quad V = \frac{v}{t}, \quad X = \frac{x}{a}, \quad Y = \frac{y}{b} \quad (6)$$

$$Q = \frac{q}{E}, \quad c = \frac{a}{b}, \quad e = \frac{t}{a}, \quad X_2 = 1, \quad Y_2 = \sqrt[2k]{1 - X^{2k}} \quad (7)$$

$$M_x = \frac{a}{D} m_x, \quad M_y = \frac{a}{D} m_y, \quad \alpha_x = \frac{D}{a^3 q} M_x, \quad \alpha_y = \frac{D}{a^3 q} M_y \quad (8)$$

Hence, substituting the non-dimensional variables into Eqs. (2)-(4)

$$V_1 = \frac{Et^3}{2(1-\mu^2)c} \int_0^{x_2} \int_0^{y_2} \left\{ \frac{e^2}{12} \left( \frac{\partial^2 W}{\partial X^2} \right)^2 + \frac{c^4 e^2}{12} \left( \frac{\partial^2 W}{\partial Y^2} \right)^2 + \frac{c^2 e^2}{6} \left( \frac{\partial^2 W}{\partial X^2} \right) \left( \frac{\partial^2 W}{\partial Y^2} \right) \right. \\ \left. - \frac{(1-\mu)c^2 e^2}{6} \left[ \left( \frac{\partial^2 W}{\partial X^2} \right) \left( \frac{\partial^2 W}{\partial Y^2} \right) - \left( \frac{\partial^2 W}{\partial X \partial Y} \right)^2 \right] \right\} dX dY \quad (9)$$

$$V_2 = \frac{Et^3}{2(1-\mu^2)c} \int_0^{x_2} \int_0^{y_2} \left\{ \left[ \left( \frac{\partial U}{\partial X} \right)^2 + e \left( \frac{\partial U}{\partial X} \right) \left( \frac{\partial W}{\partial X} \right)^2 + c^2 \left( \frac{\partial V}{\partial Y} \right)^2 \right] \right. \\ \left[ + c^3 e \left( \frac{\partial V}{\partial Y} \right) \left( \frac{\partial W}{\partial Y} \right)^2 \right] \\ + \frac{e^2}{4} \left[ \left( \frac{\partial W}{\partial X} \right)^4 + c^4 \left( \frac{\partial W}{\partial Y} \right)^4 + 2c^2 \left( \frac{\partial W}{\partial X} \right)^2 \left( \frac{\partial W}{\partial Y} \right)^2 \right] \\ \left. + 2\mu \left[ c \left( \frac{\partial U}{\partial X} \right) \left( \frac{\partial V}{\partial Y} \right) + \frac{ce}{2} \left( \frac{\partial V}{\partial Y} \right) \left( \frac{\partial W}{\partial X} \right)^2 \right] \right. \\ \left. + \frac{c^2 e}{2} \left( \frac{\partial U}{\partial X} \right) \left( \frac{\partial W}{\partial Y} \right)^2 \right] \right\} dX dY$$

$$\left[ \begin{aligned} & c^2 \left( \frac{\partial U}{\partial Y} \right)^2 + 2c \left( \frac{\partial U}{\partial Y} \right) \left( \frac{\partial V}{\partial X} \right) + \left( \frac{\partial V}{\partial X} \right)^2 \\ & + \frac{(1-\mu)}{2} + 2c^2 e \left( \frac{\partial U}{\partial Y} \right) \left( \frac{\partial W}{\partial X} \right) \left( \frac{\partial W}{\partial Y} \right) \\ & + 2ce \left( \frac{\partial V}{\partial X} \right) \left( \frac{\partial W}{\partial X} \right) \left( \frac{\partial W}{\partial Y} \right) \end{aligned} \right] \quad (10)$$

$$V_3 = -\frac{Et^3}{e^2 c} Q \int_0^{x_2} \int_0^{y_2} W \, dX \, dY \quad (11)$$

$$M_x = -e \left( \frac{\partial^2 W}{\partial X^2} + \mu c^2 \frac{\partial^2 W}{\partial Y^2} \right), \quad M_y = -e \left( c^2 \frac{\partial^2 W}{\partial Y^2} + \mu \frac{\partial^2 W}{\partial X^2} \right) \quad (12)$$

are obtained. Due to the symmetry of the plate geometry, a quarter of the plate is used in the computations.

### 3. Ritz method

The plate is considered to be resting on symmetrically distributed four point supports located at the vertices of a concentric rectangle. The positions of the point supports which prevent both deflection and in-plane displacements are defined by

$$d_x = \Delta_1 a, \quad d_y = \Delta_2 b \quad (13)$$

where  $d_x$  and  $d_y$  denote the distance from the central axes of the plate (Fig. 1).

The procedure of the Ritz method can be summarized in two steps. First, an admissible solution which satisfies the Dirichlet (essential) boundary conditions and contains unknown coefficients is chosen (Kwon and Bang 2000). The admissible functions must be at least  $p$  times differentiable (Monterrubio and Ilanko 2012). Next, the functional into which the assumed solution is substituted is minimized and the unknown coefficients are found (Kwon and Bang 2000). However, if it is difficult to construct a series of assumed functions that satisfies all the prescribed boundary conditions, the Ritz method combined with Lagrange multipliers is employed (Szilard 2004). Therefore, the geometrical boundary conditions at the point support at  $(X, Y) = (\Delta_1, \Delta_2)$  are satisfied by the subsidiary conditions  $G_1, G_2, G_3$  given by (Brebbia 1984)

$$G_1 = W(\Delta_1, \Delta_2), \quad G_2 = U(\Delta_1, \Delta_2), \quad G_3 = V(\Delta_1, \Delta_2) \quad (14)$$

Consequently, the modified functional  $F$  becomes

$$F = V_1 + V_2 + V_3 + \sum_{i=1}^3 \lambda_i G_i \quad (15)$$

where  $\lambda_1, \lambda_2, \lambda_3$  are the Lagrange multipliers. The transverse deflection surface  $W$ , and the

horizontal displacements  $U$ , and  $V$  are constructed as two dimensional complete polynomials given by

$$W = W(X, Y) = \sum_{m=0}^{d/2} \sum_{n=0}^m A_{mn} X^{2n} Y^{2(m-n)}, \quad U = U(X, Y) = X \sum_{m=0}^{d/2} \sum_{n=0}^m B_{mn} X^{2n} Y^{2(m-n)} \quad (16)$$

$$V = V(X, Y) = Y \sum_{m=0}^{d/2} \sum_{n=0}^m C_{mn} X^{2n} Y^{2(m-n)}, \quad r = \frac{(d+2)(d+4)}{8} \quad (17)$$

where the unknowns  $A_{mn}$ ,  $B_{mn}$ ,  $C_{mn}$  and  $\lambda_i$  are determined from the minimum potential energy principle by (Wang *et al.* 2002)

$$\frac{\partial F}{\partial A_{mn}} = 0, \quad \frac{\partial F}{\partial B_{mn}} = 0, \quad \frac{\partial F}{\partial C_{mn}} = 0, \quad \frac{\partial F}{\partial \lambda_i} = 0, \quad i = 1, 2, 3, \quad (18)$$

Eq. (18) yields a system of nonlinear equations due to  $V_2$ . The components of linear and nonlinear coefficient matrices  $C_L$ , and  $C_N$  are introduced by

$$[C_L] = \begin{bmatrix} \left[ \frac{\partial^2 F_L}{\partial A_{mn}^2} \right]_{rxr} & \left[ \frac{\partial^2 F_L}{\partial A_{mn} \partial B_{mn}} \right]_{rxr} & \left[ \frac{\partial^2 F_L}{\partial A_{mn} \partial C_{mn}} \right]_{rxr} & \left[ \frac{\partial^2 F_L}{\partial A_{mn} \partial \lambda_i} \right]_{rx3} \\ \left[ \frac{\partial^2 F_L}{\partial B_{mn} \partial A_{mn}} \right]_{rxr} & \left[ \frac{\partial^2 F_L}{\partial B_{mn}^2} \right]_{rxr} & \left[ \frac{\partial^2 F_L}{\partial B_{mn} \partial C_{mn}} \right]_{rxr} & \left[ \frac{\partial^2 F_L}{\partial B_{mn} \partial \lambda_i} \right]_{rx3} \\ \left[ \frac{\partial^2 F_L}{\partial C_{mn} \partial A_{mn}} \right]_{rxr} & \left[ \frac{\partial^2 F_L}{\partial C_{mn} \partial B_{mn}} \right]_{rxr} & \left[ \frac{\partial^2 F_L}{\partial C_{mn}^2} \right]_{rxr} & \left[ \frac{\partial^2 F_L}{\partial C_{mn} \partial \lambda_i} \right]_{rx3} \\ \left[ \frac{\partial^2 F_L}{\partial \lambda_i \partial A_{mn}} \right]_{3xr} & \left[ \frac{\partial^2 F_L}{\partial \lambda_i \partial B_{mn}} \right]_{3xr} & \left[ \frac{\partial^2 F_L}{\partial \lambda_i \partial C_{mn}} \right]_{3xr} & \left[ \frac{\partial^2 F_L}{\partial \lambda_i^2} \right]_{3 \times 3} \end{bmatrix} \quad (19)$$

$$[C_N] = \begin{bmatrix} \left[ \frac{\partial^2 F_N}{\partial A_{mn}^2} \right]_{rxr} & \left[ \frac{\partial^2 F_N}{\partial A_{mn} \partial B_{mn}} \right]_{rxr} & \left[ \frac{\partial^2 F_N}{\partial A_{mn} \partial C_{mn}} \right]_{rxr} & \left[ \frac{\partial^2 F_N}{\partial A_{mn} \partial \lambda_i} \right]_{rx3} \\ \left[ \frac{\partial^2 F_N}{\partial B_{mn} \partial A_{mn}} \right]_{rxr} & \left[ \frac{\partial^2 F_N}{\partial B_{mn}^2} \right]_{rxr} & \left[ \frac{\partial^2 F_N}{\partial B_{mn} \partial C_{mn}} \right]_{rxr} & \left[ \frac{\partial^2 F_N}{\partial B_{mn} \partial \lambda_i} \right]_{rx3} \\ \left[ \frac{\partial^2 F_N}{\partial C_{mn} \partial A_{mn}} \right]_{rxr} & \left[ \frac{\partial^2 F_N}{\partial C_{mn} \partial B_{mn}} \right]_{rxr} & \left[ \frac{\partial^2 F_N}{\partial C_{mn}^2} \right]_{rxr} & \left[ \frac{\partial^2 F_N}{\partial C_{mn} \partial \lambda_i} \right]_{rx3} \\ \left[ \frac{\partial^2 F_N}{\partial \lambda_i \partial A_{mn}} \right]_{3xr} & \left[ \frac{\partial^2 F_N}{\partial \lambda_i \partial B_{mn}} \right]_{3xr} & \left[ \frac{\partial^2 F_N}{\partial \lambda_i \partial C_{mn}} \right]_{3xr} & \left[ \frac{\partial^2 F_N}{\partial \lambda_i^2} \right]_{3 \times 3} \end{bmatrix} \quad (20)$$

where

$$F_L = \sum_{i=1}^3 \lambda_i G_i + \frac{Et^3}{2(1-\mu^2)c} \int_0^{x_2} \int_0^{y_2} \left\{ \begin{aligned} & \left[ \frac{e^2}{12} \left( \frac{\partial^2 W}{\partial X^2} \right)^2 + \frac{c^4 e^2}{12} \left( \frac{\partial^2 W}{\partial Y^2} \right)^2 + \frac{c^2 e^2}{6} \left( \frac{\partial^2 W}{\partial X^2} \right) \left( \frac{\partial^2 W}{\partial Y^2} \right) \right. \\ & - \frac{(1-\mu)c^2 e^2}{6} \left[ \left( \frac{\partial^2 W}{\partial X^2} \right) \left( \frac{\partial^2 W}{\partial Y^2} \right) - \left( \frac{\partial^2 W}{\partial X \partial Y} \right)^2 \right] \\ & + \left( \frac{\partial U}{\partial X} \right)^2 + c^2 \left( \frac{\partial V}{\partial Y} \right)^2 + 2\mu c \left( \frac{\partial U}{\partial X} \right) \left( \frac{\partial V}{\partial Y} \right) \\ & \left. + \frac{(1-\mu)}{2} \left[ c^2 \left( \frac{\partial U}{\partial Y} \right)^2 + 2c \left( \frac{\partial U}{\partial Y} \right) \left( \frac{\partial V}{\partial X} \right) + \left( \frac{\partial V}{\partial X} \right)^2 \right] \right\} dX dY \quad (21) \end{aligned} \right.$$

$$F_N = \frac{Et^3}{2(1-\mu^2)c} \int_0^{x_2} \int_0^{y_2} \left\{ \begin{aligned} & e \left( \frac{\partial U}{\partial X} \right) \left( \frac{\partial W}{\partial X} \right)^2 + c^3 e \left( \frac{\partial V}{\partial Y} \right) \left( \frac{\partial W}{\partial Y} \right)^2 \\ & + \frac{e^2}{4} \left[ \left( \frac{\partial W}{\partial X} \right)^4 + c^4 \left( \frac{\partial W}{\partial Y} \right)^4 + 2c^2 \left( \frac{\partial W}{\partial X} \right)^2 \left( \frac{\partial W}{\partial Y} \right)^2 \right] \\ & + 2\mu \left[ \frac{ce}{2} \left( \frac{\partial V}{\partial Y} \right) \left( \frac{\partial W}{\partial X} \right)^2 + \frac{c^2 e}{2} \left( \frac{\partial U}{\partial X} \right) \left( \frac{\partial W}{\partial Y} \right)^2 \right] \\ & + \frac{(1-\mu)}{2} \left[ 2c^2 e \left( \frac{\partial U}{\partial Y} \right) \left( \frac{\partial W}{\partial X} \right) \left( \frac{\partial W}{\partial Y} \right) \right. \\ & \left. + 2ce \left( \frac{\partial V}{\partial X} \right) \left( \frac{\partial W}{\partial X} \right) \left( \frac{\partial W}{\partial Y} \right) \right] \end{aligned} \right\} dX dY \quad (22)$$

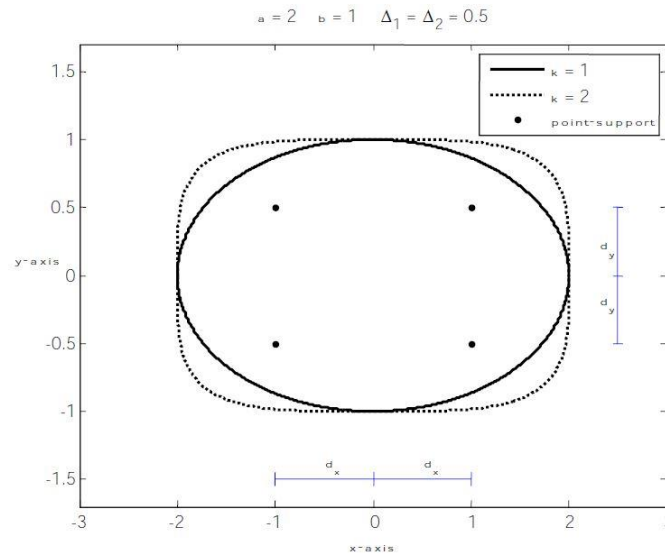


Fig. 1 Geometry of the plate and the location of the point supports

#### 4. Newton-Raphson method

The Newton-Raphson (NR) method is one of the best known iterative techniques in finding the roots of nonlinear equations. The algorithm requires the evaluation of two functions per iteration: the Jacobian matrix which involves the partial derivatives, and the function vector. Depending on the initial guess the NR method may converge fast to the solution. After each iteration, convergence is checked, and the values of the unknowns are updated. The iterations stop when convergence is obtained, and the error is less than the desired tolerance value provided that the maximum number of iterations is not exceeded (Mathews 1992).

In the current paper totally  $3r+3$  unknowns were considered, and the zero vector (i.e., the linear solution) was assumed to be the initial guess. For each iteration the Euclidian norm of the function vector was determined, and  $E_a$ , and  $E_r$  were computed. When  $E_a \leq \varepsilon$ , or  $|E_r| \leq \varepsilon$ , the iterations were stopped.

#### 5. Numerical results

The integration with respect to “ $X$ ” was evaluated numerically by the 12-point Gaussian quadrature technique (Maron and Lopez 1991). The analysis was made for  $0.1 \leq \Delta \leq 0.7$  where  $\Delta = \Delta_1 = \Delta_2$ . Unless otherwise stated the semi-minor axis, the nondimensional load, and the Poisson’s ratio were considered to be  $b=1\text{m}$ ,  $Q=5 \times 10^{-10}$ , and  $\mu=0.3$ . As a special case, a corner supported plate was examined for  $\Delta = 1/\sqrt[2k]{2}$ . Convergence studies were performed and  $d=8$  was found to be sufficient for admissible accuracy. The results of the central deflection computed in the present study for varying support location were compared with those of elliptical ( $k=1$ ) and rectangular ( $k=250$ ) plates, and good agreement was obtained (Table 1). The location of the point where the maximum deflection develops was highlighted by comparing the deflections at the points  $R_j(X, Y)$  where  $j=0,1,2,3$ . The coordinates of these points are  $R_0(0,0)$ ,  $R_1(1,0)$ ,  $R_2(0,1)$ ,  $R_3(1/\sqrt[2k]{2}, 1/\sqrt[2k]{2})$ . Therefore, the deflections  $\alpha_0, \alpha_1, \alpha_2, \alpha_3$  where the subscripts denote the points  $R_0, R_1, R_2, R_3$ , respectively, were computed for varying support location (Table 2). The central deflection of a

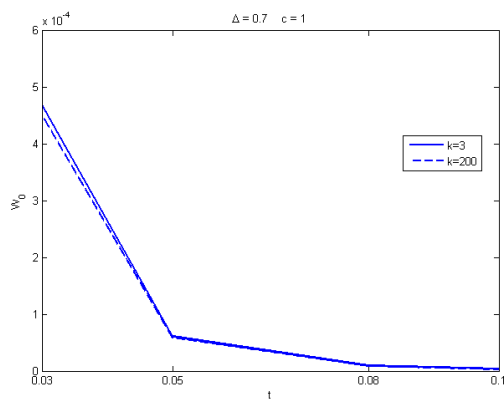


Fig. 2 Influence of the thickness on the central deflection ( $c=1$ )

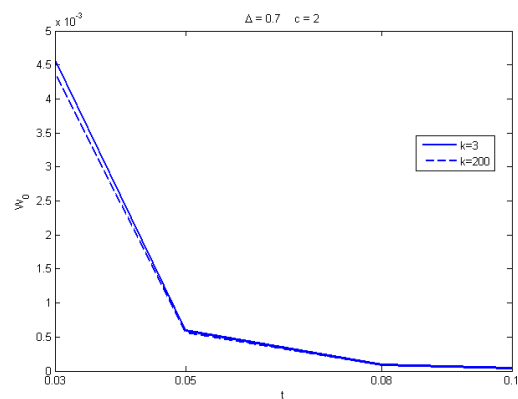


Fig. 3 Influence of the thickness on the central deflection ( $c=2$ )

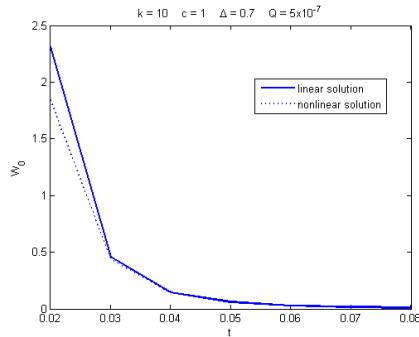


Fig. 4 Comparison of the linear and nonlinear solutions ( $c=1$ ,  $Q=5 \times 10^{-7}$ )

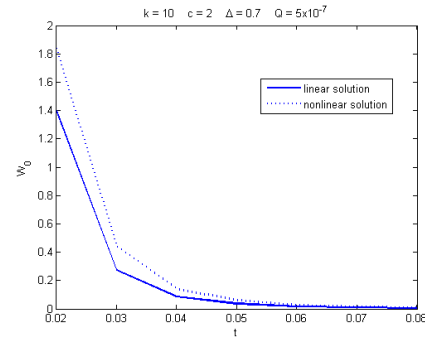


Fig. 5 Comparison of the linear and nonlinear solutions ( $c=2$ ,  $Q=5 \times 10^{-7}$ )

Table 1 Convergence study and comparison of the central deflection results for a point-supported super-elliptical plate under uniform pressure ( $t=0.01$  m)

$k$	$c$	$\Delta$	$\mu$	$d$	$a_0$	Reference	Geometry of the plate
1	1	$1/\sqrt[2k]{2}$	0.17	4	0.0816	(Williams and Brinson 1974) (Altekin 2010)	Circular plate
				6	0.0835		
				8	0.0849		
					0.0861		
					0.0858		
1	1	$0.9/\sqrt{2}$	0.17	4	0.0456	(Williams and Brinson 1974) (Altekin 2010)	Circular plate
				6	0.0490		
				8	0.0491		
					0.0500		
					0.0499		
250	1	$1/\sqrt[2k]{2}$	0.28	4	0.4060	(Shanmugam <i>et al.</i> 1988) (Altekin 2010)	Square plate
				6	0.4053		
				8	0.4055		
					0.4052		
					0.4060		
250	1	$1/\sqrt[2k]{2}$	0.30	4	0.4052	(Wang <i>et al.</i> 2002)* (Rajaiah and Rao 1978) (Altekin 2010)	Square plate
				6	0.4045		
				8	0.4048		
					0.4096		
					0.4080		
1	2	$1/\sqrt[2k]{2}$	0.30	4	0.0475	(Altekin and Altay 2008)	Elliptical plate
				6	0.0491		
				8	0.0495		
					0.0497		
					0.0497		

\*The results obtained from Wang *et al.* (2002), Shanmugam *et al.* (1988), and Szilard (1974) were scaled by the author.



Table 1 Continued

20	2	$1/\sqrt[2k]{2}$	0.30	4	0.2123	Super-elliptical plate (Altekin and Altay 2008)
				6	0.2132	
				8	0.2134	
					0.2140	
250	2	$1/\sqrt[2k]{2}$	1/6	4	0.2227	Rectangular plate (Szilard 1974)
				6	0.2238	
				8	0.2239	
					0.2240	

Table 2 Location of  $\alpha_{\max}$  for varying support location ( $t=0.05, \mu=0.3$ )

$c$	$\Delta$	$\alpha_0$	$\alpha_1$	$\alpha_2$	$\alpha_3$
1	$\Delta=0.1$				$\alpha_{\max}=\alpha_3$
	$0.2 \leq \Delta \leq 0.4$		$\alpha_{\max}=\alpha_1$ for $k=1$		$\alpha_{\max}=\alpha_3$ for $k \geq 2$
	$\Delta=0.5$		$\alpha_{\max}=\alpha_0$ for $k \leq 2$		$\alpha_{\max}=\alpha_3$ for $k \geq 3$
	$0.6 \leq \Delta \leq 0.7$	$\alpha_{\max}=\alpha_0$			
2	$0.1 \leq \Delta \leq 0.4$		$\alpha_{\max}=\alpha_1$		
	$\Delta=0.5$		$\alpha_{\max}=\alpha_1$ for $k \geq 2$	$\alpha_{\max}=\alpha_2$ for $k=1$	
	$0.6 \leq \Delta \leq 0.7$			$\alpha_{\max}=\alpha_2$	

Table 3 Central deflection  $\alpha_0$  for varying support location ( $c=1, t=0.05, \mu=0.3$ )

$2k$	$\Delta=0.1$	$\Delta=0.2$	$\Delta=0.3$	$\Delta=0.4$	$\Delta=0.5$	$\Delta=0.6$	$\Delta=0.7$
2	-0.0035	-0.0096	-0.0111	-0.0054	0.0087	0.0341	0.0780
4	-0.0042	-0.0123	-0.0162	-0.0116	0.0023	0.0282	0.0716
6	-0.0044	-0.0131	-0.0178	-0.0137	0.0001	0.0260	0.0696
8	-0.0044	-0.0134	-0.0184	-0.0146	-0.0009	0.0250	0.0687
10	-0.0045	-0.0135	-0.0188	-0.0151	-0.0015	0.0244	0.0681
12	-0.0045	-0.0136	-0.0190	-0.0154	-0.0018	0.0241	0.0678
14	-0.0045	-0.0136	-0.0191	-0.0156	-0.0020	0.0239	0.0676
16	-0.0045	-0.0137	-0.0192	-0.0158	-0.0021	0.0238	0.0675
18	-0.0045	-0.0137	-0.0193	-0.0159	-0.0022	0.0237	0.0674
20	-0.0045	-0.0137	-0.0193	-0.0159	-0.0023	0.0236	0.0673
22	-0.0045	-0.0137	-0.0193	-0.0160	-0.0024	0.0235	0.0672
24	-0.0045	-0.0137	-0.0194	-0.0160	-0.0024	0.0235	0.0672
26	-0.0045	-0.0137	-0.0194	-0.0161	-0.0024	0.0235	0.0672
28	-0.0045	-0.0137	-0.0194	-0.0161	-0.0025	0.0234	0.0671
30	-0.0045	-0.0137	-0.0194	-0.0161	-0.0025	0.0234	0.0671
32	-0.0045	-0.0137	-0.0194	-0.0161	-0.0025	0.0234	0.0671
34	-0.0045	-0.0137	-0.0194	-0.0161	-0.0025	0.0234	0.0671
36	-0.0045	-0.0137	-0.0195	-0.0162	-0.0025	0.0234	0.0671
38	-0.0045	-0.0137	-0.0195	-0.0162	-0.0025	0.0234	0.0671
40	-0.0045	-0.0137	-0.0195	-0.0162	-0.0026	0.0233	0.0670
100	-0.0045	-0.0137	-0.0195	-0.0163	-0.0026	0.0233	0.0670
400	-0.0045	-0.0137	-0.0195	-0.0163	-0.0027	0.0233	0.0670

Table 4 Central deflection  $\alpha_0$  for varying support location ( $c=2, t=0.05, \mu=0.3$ )

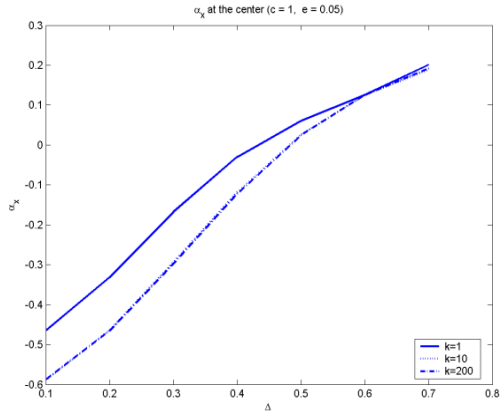
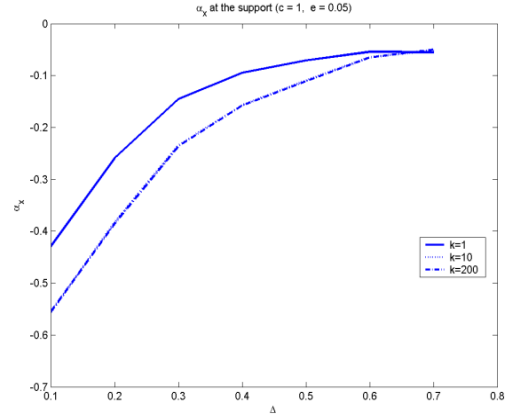
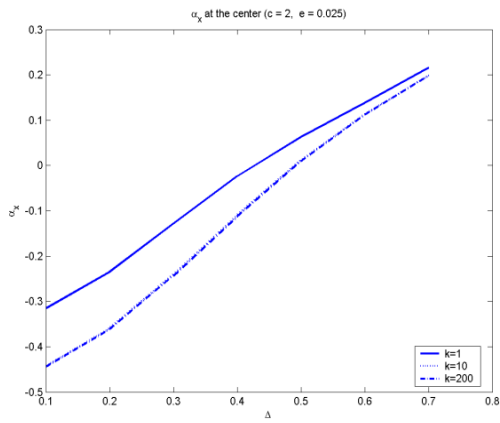
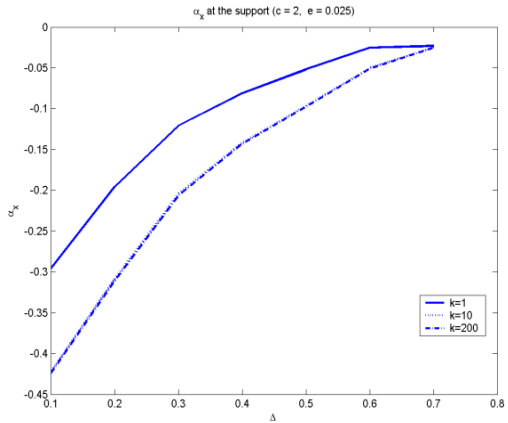
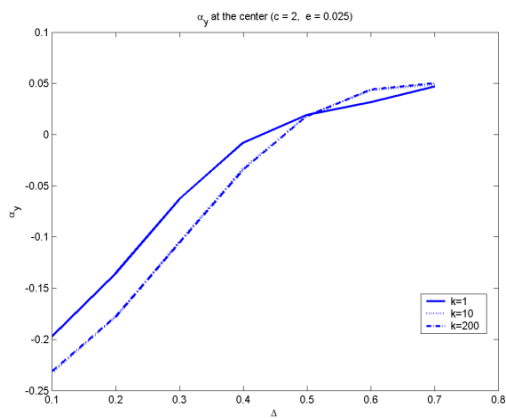
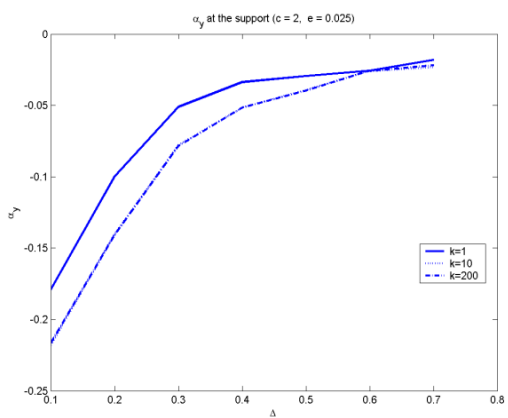
$2k$	$\Delta=0.1$	$\Delta=0.2$	$\Delta=0.3$	$\Delta=0.4$	$\Delta=0.5$	$\Delta=0.6$	$\Delta=0.7$
2	-0.0015	-0.0045	-0.0057	-0.0030	0.0050	0.0206	0.0471
4	-0.0019	-0.0060	-0.0086	-0.0068	0.0008	0.0167	0.0437
6	-0.0021	-0.0065	-0.0095	-0.0082	-0.0007	0.0152	0.0424
8	-0.0021	-0.0067	-0.0099	-0.0088	-0.0014	0.0145	0.0418
10	-0.0021	-0.0068	-0.0101	-0.0091	-0.0017	0.0141	0.0414
12	-0.0022	-0.0068	-0.0103	-0.0093	-0.0019	0.0139	0.0412
14	-0.0022	-0.0069	-0.0103	-0.0094	-0.0021	0.0137	0.0411
16	-0.0022	-0.0069	-0.0104	-0.0095	-0.0022	0.0136	0.0410
18	-0.0022	-0.0069	-0.0104	-0.0095	-0.0022	0.0135	0.0409
20	-0.0022	-0.0069	-0.0104	-0.0096	-0.0023	0.0135	0.0408
22	-0.0022	-0.0069	-0.0105	-0.0096	-0.0023	0.0134	0.0408
24	-0.0022	-0.0069	-0.0105	-0.0096	-0.0023	0.0134	0.0408
26	-0.0022	-0.0069	-0.0105	-0.0096	-0.0024	0.0134	0.0407
28	-0.0022	-0.0069	-0.0105	-0.0096	-0.0024	0.0134	0.0407
30	-0.0022	-0.0069	-0.0105	-0.0097	-0.0024	0.0134	0.0407
32	-0.0022	-0.0069	-0.0105	-0.0097	-0.0024	0.0133	0.0407
34	-0.0022	-0.0069	-0.0105	-0.0097	-0.0024	0.0133	0.0407
36	-0.0022	-0.0069	-0.0105	-0.0097	-0.0024	0.0133	0.0407
38	-0.0022	-0.0070	-0.0105	-0.0097	-0.0024	0.0133	0.0407
40	-0.0022	-0.0070	-0.0105	-0.0097	-0.0024	0.0133	0.0407
100	-0.0022	-0.0070	-0.0106	-0.0097	-0.0025	0.0133	0.0406
400	-0.0022	-0.0070	-0.0106	-0.0097	-0.0025	0.0132	0.0406

Table 5 Coefficients determined by the method of least squares ( $c=1, t=0.05$ )

$\Delta$	$h_1$	$h_2$	$h_3$	$h_4$
0.1	-0.0041263	$-1.8096 \times 10^{-5}$	$1.8797 \times 10^{-7}$	$-3.6212 \times 10^{-10}$
0.2	-0.012063	$-8.5564 \times 10^{-5}$	$8.7633 \times 10^{-7}$	$-1.6823 \times 10^{-9}$
0.3	-0.015762	-0.00018558	$1.8827 \times 10^{-6}$	$-3.6056 \times 10^{-9}$
0.4	-0.011129	-0.00025066	$2.5317 \times 10^{-6}$	$-4.8431 \times 10^{-9}$
0.5	0.0027313	-0.00026279	$2.6557 \times 10^{-6}$	$-5.0811 \times 10^{-9}$
0.6	0.028503	-0.00025617	$2.5902 \times 10^{-6}$	$-4.9563 \times 10^{-9}$
0.7	0.072139	-0.00025284	$2.5559 \times 10^{-6}$	$-4.8903 \times 10^{-9}$

Table 6 Coefficients determined by the method of least squares ( $c=2, t=0.05$ )

$\Delta$	$h_1$	$h_2$	$h_3$	$h_4$
0.1	-0.0019079	$-1.3608 \times 10^{-5}$	$1.3914 \times 10^{-7}$	$-2.6701 \times 10^{-10}$
0.2	-0.0058758	$-5.4293 \times 10^{-5}$	$5.5282 \times 10^{-7}$	$-1.0597 \times 10^{-9}$
0.3	-0.0083431	-0.00010988	$1.1145 \times 10^{-6}$	$-2.1342 \times 10^{-9}$
0.4	-0.0065945	-0.00015446	$1.563 \times 10^{-6}$	$-2.9913 \times 10^{-9}$
0.5	0.0010771	-0.00017536	$1.7732 \times 10^{-6}$	$-3.393 \times 10^{-9}$
0.6	0.016854	-0.00017587	$1.7759 \times 10^{-6}$	$-3.3969 \times 10^{-9}$
0.7	0.043857	-0.00015756	$1.587 \times 10^{-6}$	$-3.0338 \times 10^{-9}$

Fig. 6  $\alpha_x$  at the center ( $c=1$ ,  $Q=5 \times 10^{-10}$ )Fig. 7  $\alpha_x$  at the support ( $c=1$ ,  $Q=5 \times 10^{-10}$ )Fig. 8  $\alpha_x$  at the center ( $c=2$ ,  $Q=5 \times 10^{-10}$ )Fig. 9  $\alpha_x$  at the support ( $c=2$ ,  $Q=5 \times 10^{-10}$ )Fig. 10  $\alpha_y$  at the center ( $c=2$ ,  $Q=5 \times 10^{-10}$ )Fig. 11  $\alpha_y$  at the support ( $c=2$ ,  $Q=5 \times 10^{-10}$ )

super-elliptical plate undergoing large deflection was presented for future reference (Tables 3-4). The influence of the thickness on the central deflection was investigated for four different values of  $t$  as shown on the horizontal axes of Figs. 2-3. The contribution of the nonlinearity on the central deflection was highlighted for  $0.02 \text{ m} \leq t \leq 0.08 \text{ m}$  (Figs. 4-5). The influence of the support location on the bending moment was examined for  $0.1 \leq \Delta \leq 0.7$  (Figs. 6-11).

## 6. Conclusions

It can be deduced that the perimeter of the plate is more critical than the center in terms of deflection, and the central deflection becomes dominant if  $\Delta > 0.5$  and  $c=1$  (Table 2). The results reveal that  $\Delta \cong 0.5$  yields the minimum absolute central deflection for  $k > 1$  (Tables 3-4). The influence of the super-elliptical power on the central deflection tends to weaken for  $k > 10$ . For practical considerations the central deflection computed for  $k=15$  for varying support location may be assumed to be the central deflection of a rectangular plate (Tables 3-4). The central deflection of a super-elliptical plate subject to fully applied uniform transverse pressure can be estimated approximately by the relationship given by

$$\alpha_0 = h_1 + h_2\beta + h_3\beta^2 + h_4\beta^3, \quad \beta = 2k \quad (23)$$

where  $h_1, h_2, h_3, h_4$  are the scalar coefficients which may be determined by the method of least squares (Tables 5-6).

The relation between the thickness and the central deflection is observed to be nonlinear. The influence of the super-elliptical power on the central deflection decreases if the thickness is raised (Figs. 2-3). The nonlinearity becomes significant for  $W_0 > 0.5$  for  $c=1$ , and  $W_0 > 0.4$  for  $c=2$  (Figs. 4-5). The nonlinearity is observed to be sensitive to the aspect ratio.

As  $\Delta$  is raised, the influence of  $k$  on  $M_x$  to be developed at the center and at the support decreases (Figs. 6-11).  $\Delta \cong 0.45$ , and  $\Delta \cong 0.67$  can be considered to be the optimum positions of the point support which minimize the bending moments at the supports and at the center of the plate (Figs. 6-11). The influence of the super-elliptical power on the bending moments  $\alpha_x$ , and  $\alpha_y$  is negligible for (i)  $\Delta \cong 0.6$  (Fig. 6), and  $\Delta \cong 0.5$  (Fig. 10) at the center of the plate and (ii)  $\Delta \cong 0.66$  (Fig. 7), and  $\Delta \cong 0.6$  at the support (Fig. 11). As  $\Delta$  is raised, the influence of  $k$  on the bending moments weakens (Figs. 8-9).

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## Nomenclature

$D, E, F$	flexural rigidity, Young’s modulus, modified functional
$Q, U, V, W$	nondimensional uniform pressure, nondimensional horizontal displacement along $x$ -axis, nondimensional horizontal displacement along $y$ -axis, nondimensional deflection
$A_{mn}, B_{mn}, C_{mn}$	coefficients of the shape functions
$C_L, C_N$	linear coefficient matrix, nonlinear coefficient matrix
$E_a, E_r$	absolute error, relative error
$G_1, G_2, G_3$	subsidiary conditions
$M_x, M_y$	nondimensional bending moments per unit length of the plate
$R_0, R_1, R_2, R_3$	the critical points where the maximum deflection may develop
$V_1, V_2, V_3$	strain energy of the plate due to bending, strain energy of the plate due to strain of the middle surface, potential energy of the lateral load
$a, b, c$	semi-major axis, semi-minor axis, aspect ratio
$d$	degree of the complete two-dimensional polynomial function
$p$	order of the highest differential operator in the functional
$e, k, q$	parameter of thickness, super-elliptical power, uniform pressure
$r, t$	number of terms in the shape function, thickness of the undeformed plate
$u, v, w$	horizontal displacement along $x$ -axis, horizontal displacement along $y$ -axis, deflection
$d_x, d_y$	half span along $x$ -axis, half span along $y$ -axis
$m_x, m_y$	bending moments per unit length of the plate

$\varepsilon, \mu$	tolerance value, Poisson's ratio
$\alpha_j, \lambda_i$	nondimensional deflection ( $j=0, 1, 2, 3$ ), Lagrange multiplier ( $i=1, 2, 3$ )
$\Delta_1, \Delta_2$	nondimensional half span parallel to x-axis, nondimensional half span parallel to y-axis